

## Quasi-geostrophic Energy Equation

The quasi-geostrophic energy equation can be derived from the First Law of Thermodynamics,

$$\frac{\partial T}{\partial t} + \vec{V} \cdot \nabla_p T + \omega T \frac{\partial \ln \theta}{\partial p} = \frac{1}{c_p} \frac{dQ}{dt}.$$

Replace T from the hydrostatic equation,

$$T = -\frac{p}{R} \frac{\partial \phi}{\partial p},$$

so,

$$\left( \frac{\partial}{\partial t} + \vec{V} \cdot \nabla_p + \omega \frac{\partial \ln \theta}{\partial p} \right) \left( -\frac{p}{R} \frac{\partial \phi}{\partial p} \right) = \frac{1}{c_p} \frac{dQ}{dt}.$$

Since p is constant on an isobaric surface,

$$\left( \frac{\partial}{\partial t} + \vec{V} \cdot \nabla_p + \omega \frac{\partial \ln \theta}{\partial p} \right) \left( -\frac{\partial \phi}{\partial p} \right) = \frac{R}{p c_p} \frac{dQ}{dt}.$$

Introduce the static stability parameter,  $\sigma$ ,

$$\sigma = \frac{\partial \phi}{\partial p} \frac{\partial \ln \theta}{\partial p},$$

$$\left( \frac{\partial}{\partial t} + \vec{V} \cdot \nabla_p \right) \left( -\frac{\partial \phi}{\partial p} \right) - \omega \sigma = \frac{R}{p c_p} \frac{dQ}{dt}.$$

Assume the flow is adiabatic, and that horizontal advection is done by the geostrophic wind,

$$\frac{\partial}{\partial t} \left( -\frac{\partial \phi}{\partial p} \right) + (\vec{V}_g \cdot \nabla_p) \left( -\frac{\partial \phi}{\partial p} \right) - \omega \sigma = 0.$$

Or finally,

$$\boxed{\frac{\partial}{\partial t} \left( -\frac{\partial \phi}{\partial p} \right) = -(\vec{V}_g \cdot \nabla_p) \left( -\frac{\partial \phi}{\partial p} \right) + \omega \sigma}.$$

Keeping the hydrostatic equation in mind (the quantity inside the parenthesis is proportional to temperature), the temperature of a layer can be changed by either horizontal temperature advection or vertical motion. Warm air advection and subsidence tend to increase temperature.